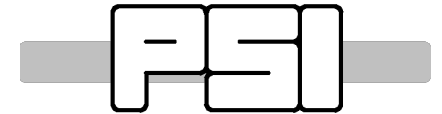
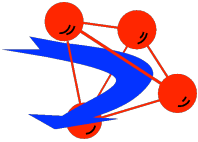


Determination of long range antiferromagnetic order by powder neutron diffraction

Vladimir Pomjakushin

Laboratory for Neutron Scattering and Imaging, LNS, Paul Scherrer Institute



During the practicum we will try to reproduce the one of the neutron diffraction experiments performed during 1946-1951 for which C.G. Shull was honoured with the Nobel Prize in 1994. We will perform neutron diffraction experiment with MnS using powder diffractometer HRPT/SINQ. From the analysis of the nuclear and magnetic Bragg peak intensities and positions we will verify the crystal and magnetic structures of manganese sulfide MnS and determine the size of the magnetic moment on manganese.

Determination of long range antiferromagnetic order by powder neutron diffraction

- Historical introduction (4)
- Some experimental details of neutron diffraction experiment at HRPT/SINQ (10)
- Reminder on nuclear and magnetic neutron structure factors (6)
- Practicum problems (6)

1994 Nobel Prize in Physics

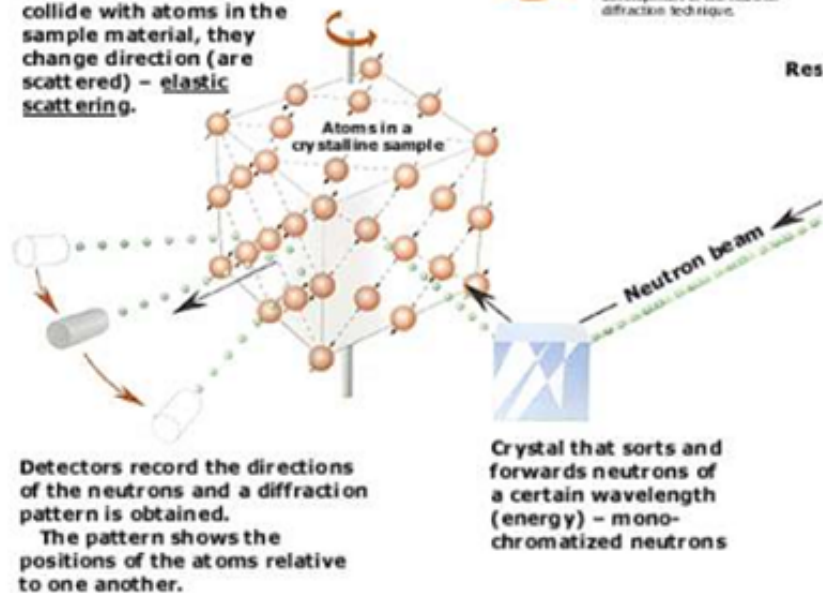
Clifford G. Shull
1915 – 2001, USA



Clifford G. Shull, MIT, Cambridge, Massachusetts, USA, receives one half of the 1994 Nobel Prize in Physics for development of the neutron diffraction technique.

Neutrons show where atoms are

When the neutrons collide with atoms in the sample material, they change direction (are scattered) – elastic scattering.



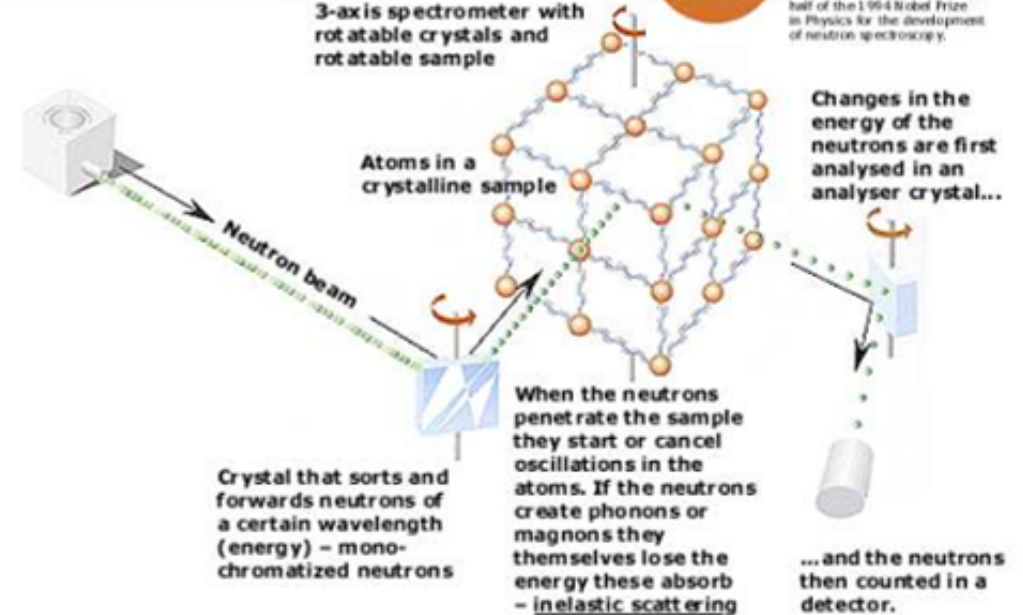
Bertram N. Brockhouse
1918 – 2003, Canada



Bertram N. Brockhouse, McMaster University, Hamilton, Ontario, Canada, receives one half of the 1994 Nobel Prize in Physics for the development of neutron spectroscopy.

Neutrons show what atoms do

3-axis spectrometer with rotatable crystals and rotatable sample



Neutron Diffraction by Paramagnetic and Antiferromagnetic Substances

C. G. SHULL, W. A. STRAUSSER, AND E. O. WOLLAN
Oak Ridge National Laboratory, Oak Ridge, Tennessee

(Received March 2, 1951)

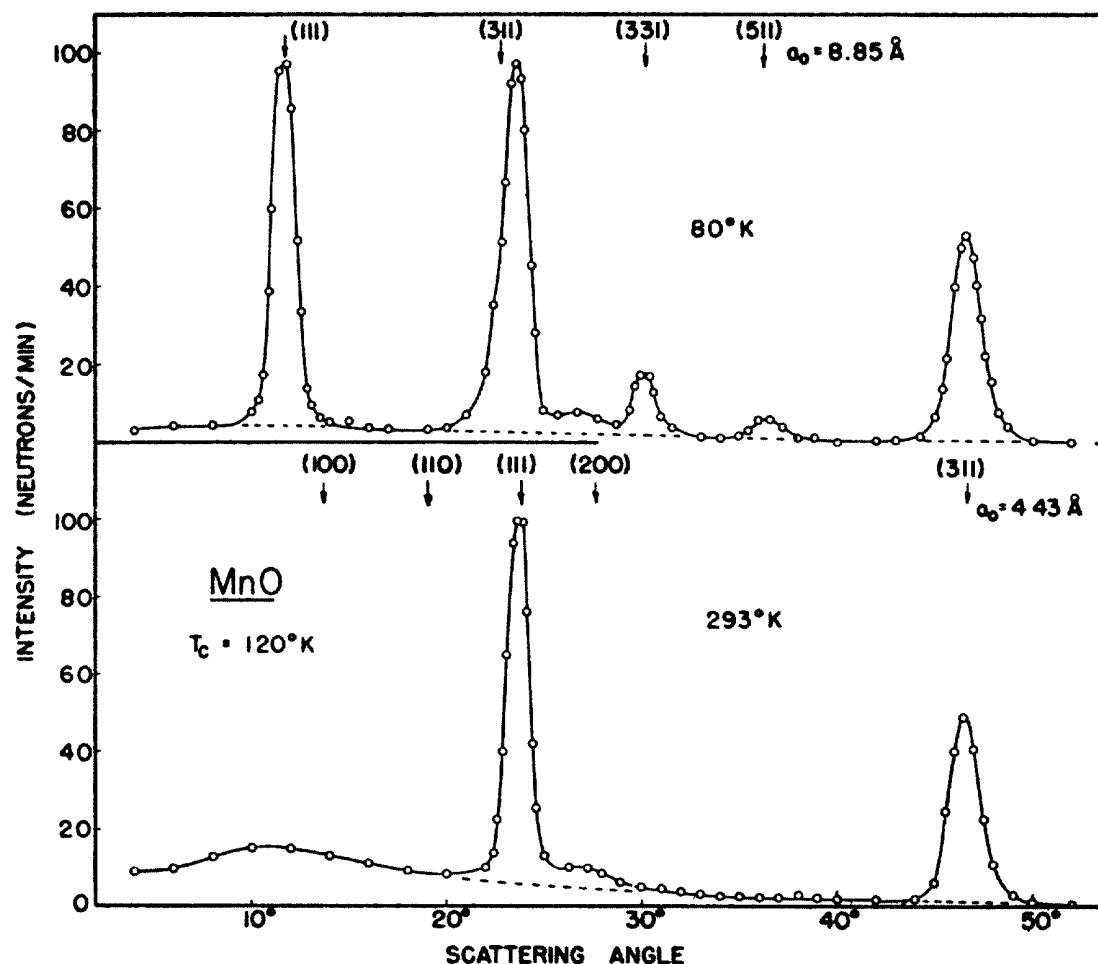


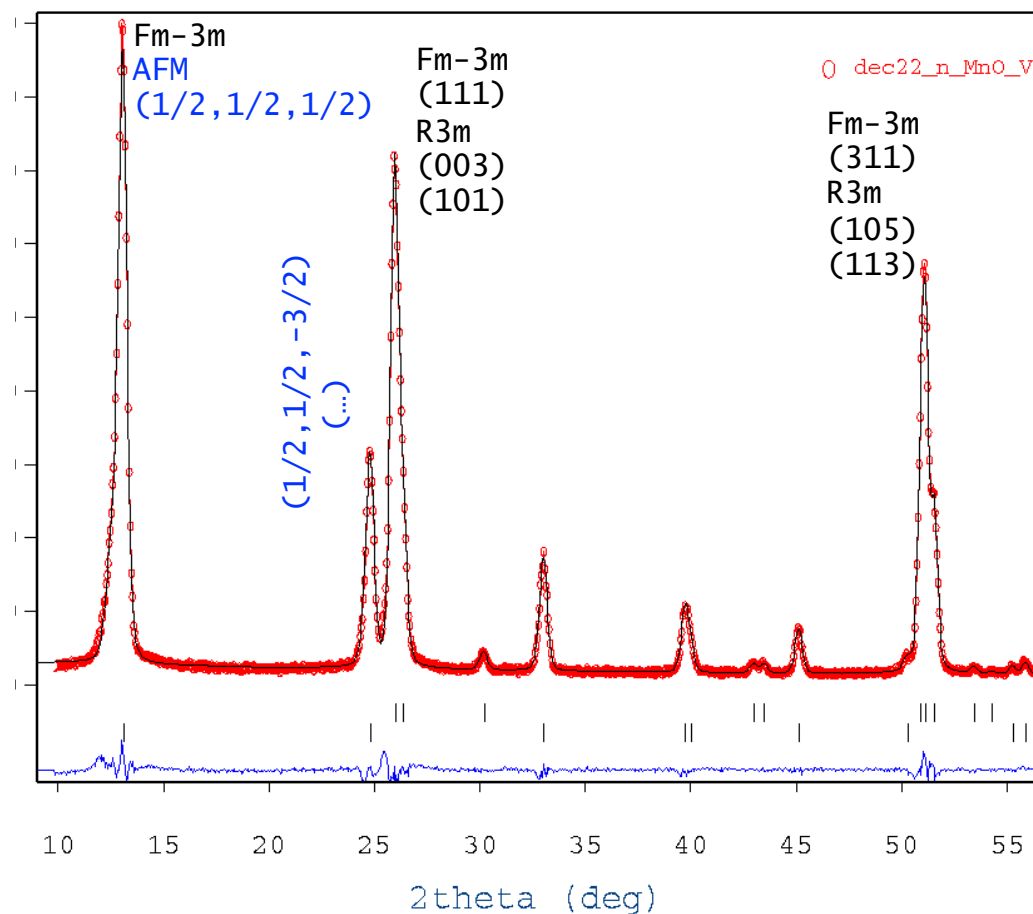
FIG. 4. Neutron diffraction patterns for MnO taken at liquid nitrogen and room temperatures. The patterns have been corrected for the various forms of extraneous, diffuse scattering mentioned in the text. Four extra antiferromagnetic reflections are to be noticed in the low temperature pattern.

HRPT/SINQ nowadays

$\lambda=1.15\text{\AA}$, MnO @ 2K.

Rhombohedral distortions are explicitly seen

R-3m and $k=003/2$



reaction by Paramagnetic and Antiferromagnetic
 $\lambda=1.057\text{\AA}$

C. G. SHULL, W. A. STRAUSSER, AND E. O. WOLLAN
Oak Ridge National Laboratory, Oak Ridge, Tennessee

(Received March 2, 1951)

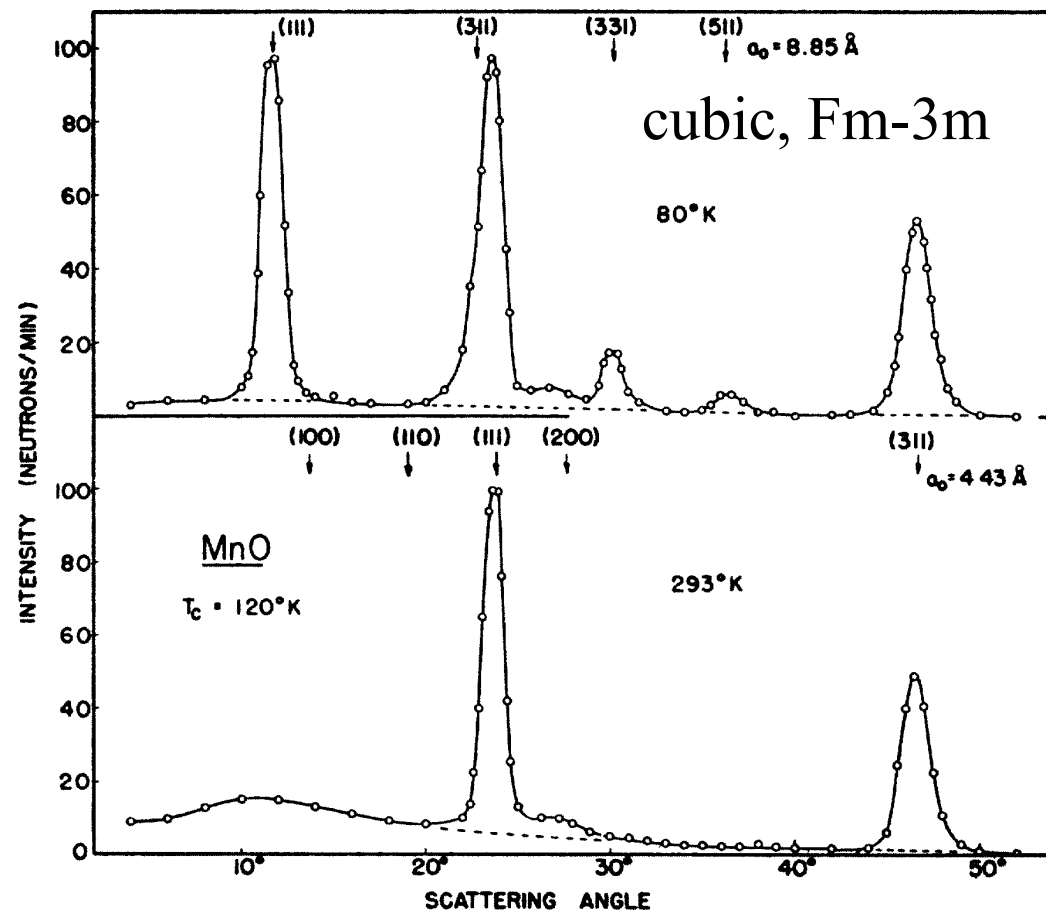
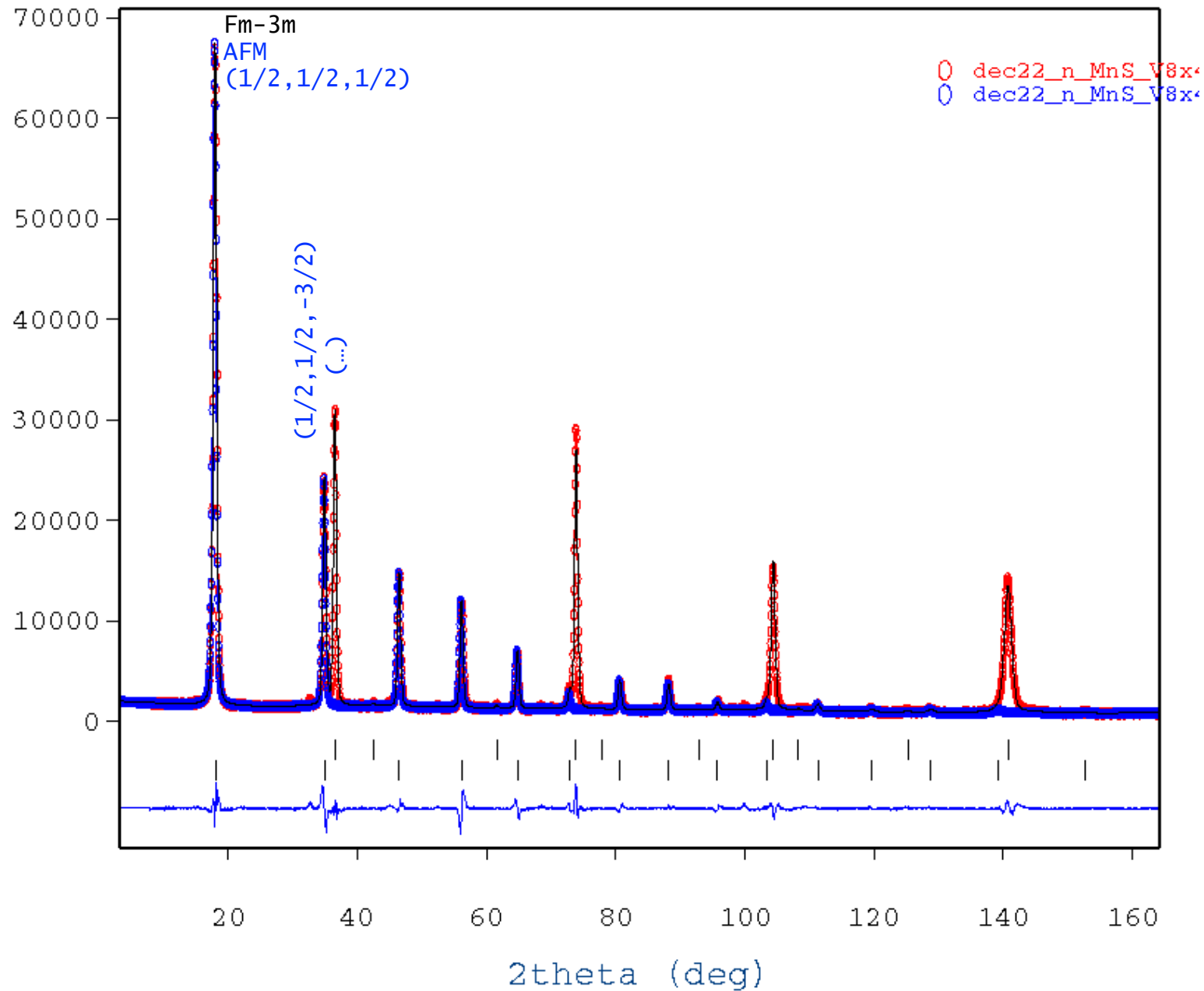


FIG. 4. Neutron diffraction patterns for MnO taken at liquid nitrogen and room temperatures. The patterns have been corrected for the various forms of extraneous, diffuse scattering mentioned in the text. Four extra antiferromagnetic reflections are to be noticed in the low temperature pattern.

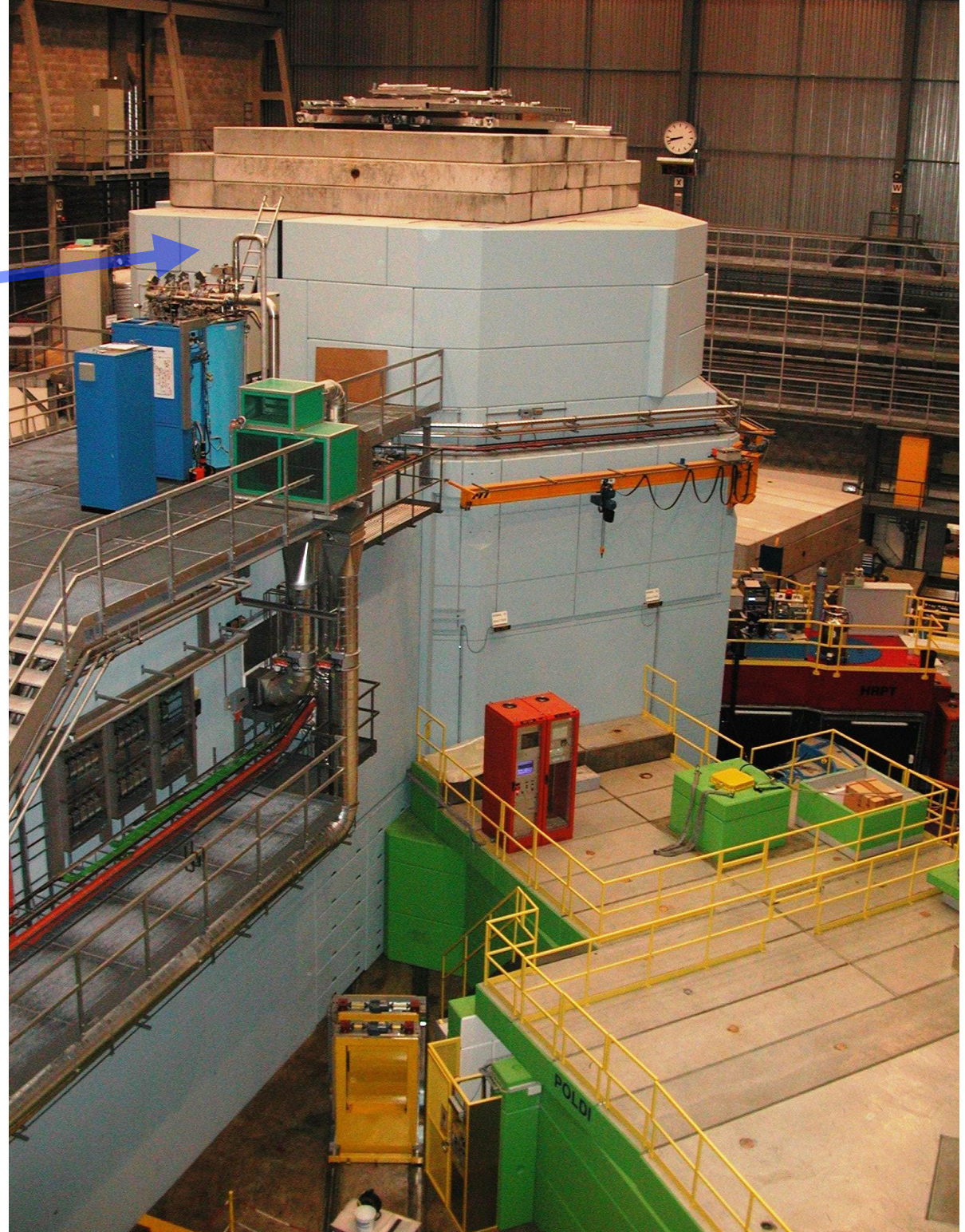
HRPT, 1.9Å, MnS @ T=2K, pseudo-cubic Fm-3m



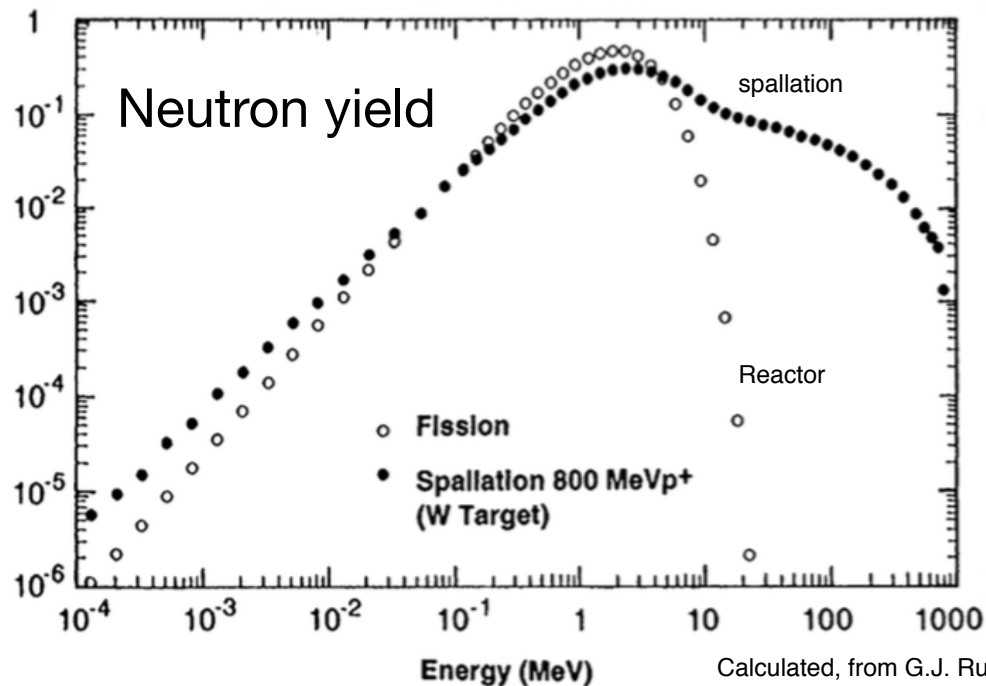
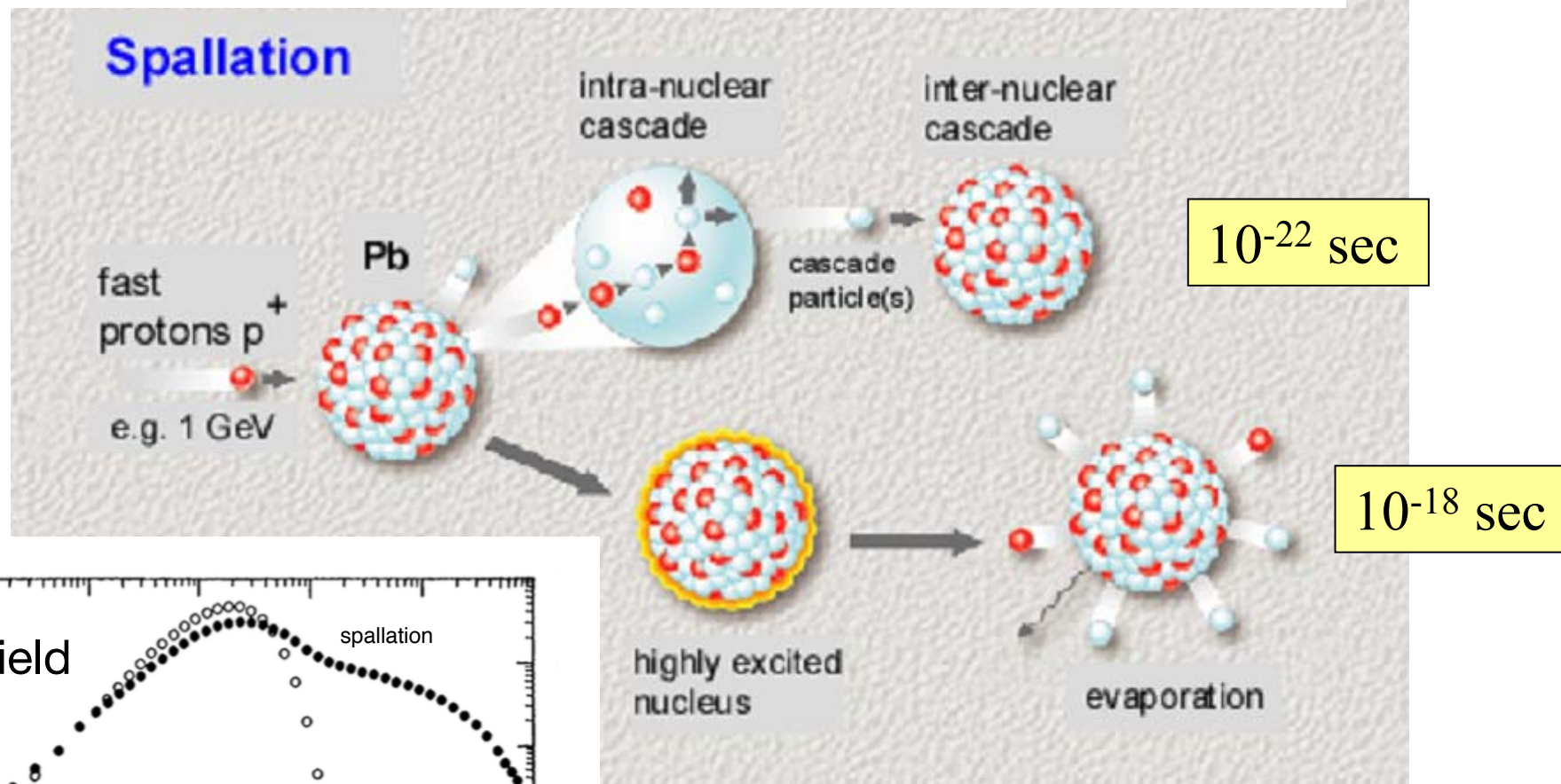
**Some experimental details of neutron
diffraction experiment at HRPT/SINQ**

SINQ hall

The spallation neutron source **SINQ** is a continuous source - the first and the only of its kind in the world - with a **flux of about 10^{14} n/cm²/s**. Beside thermal neutrons, a cold moderator of liquid deuterium (cold source) slows neutrons down and shifts their spectrum to lower energies.

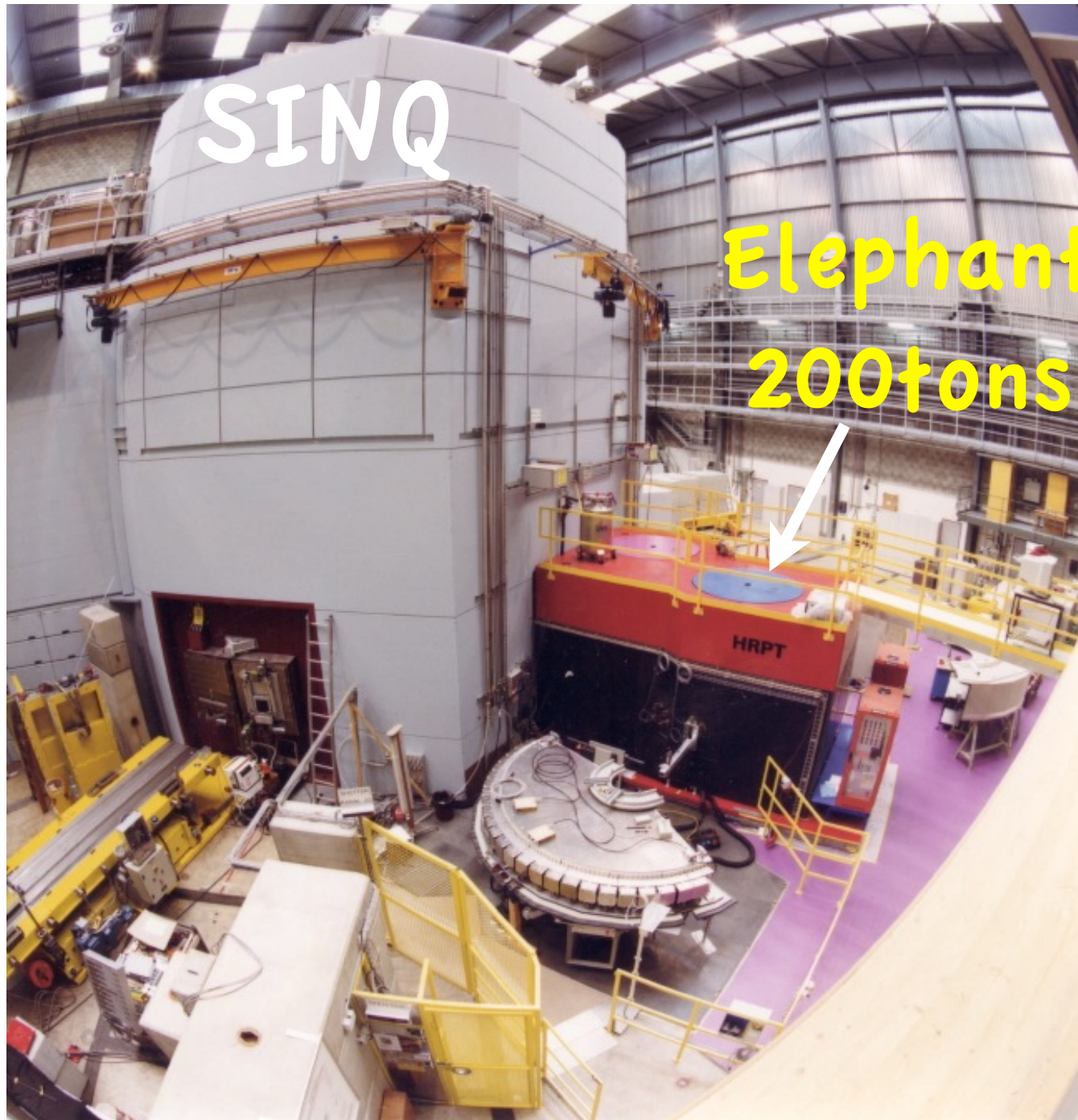


Spallation, SINQ 590MeV protons, Pb target



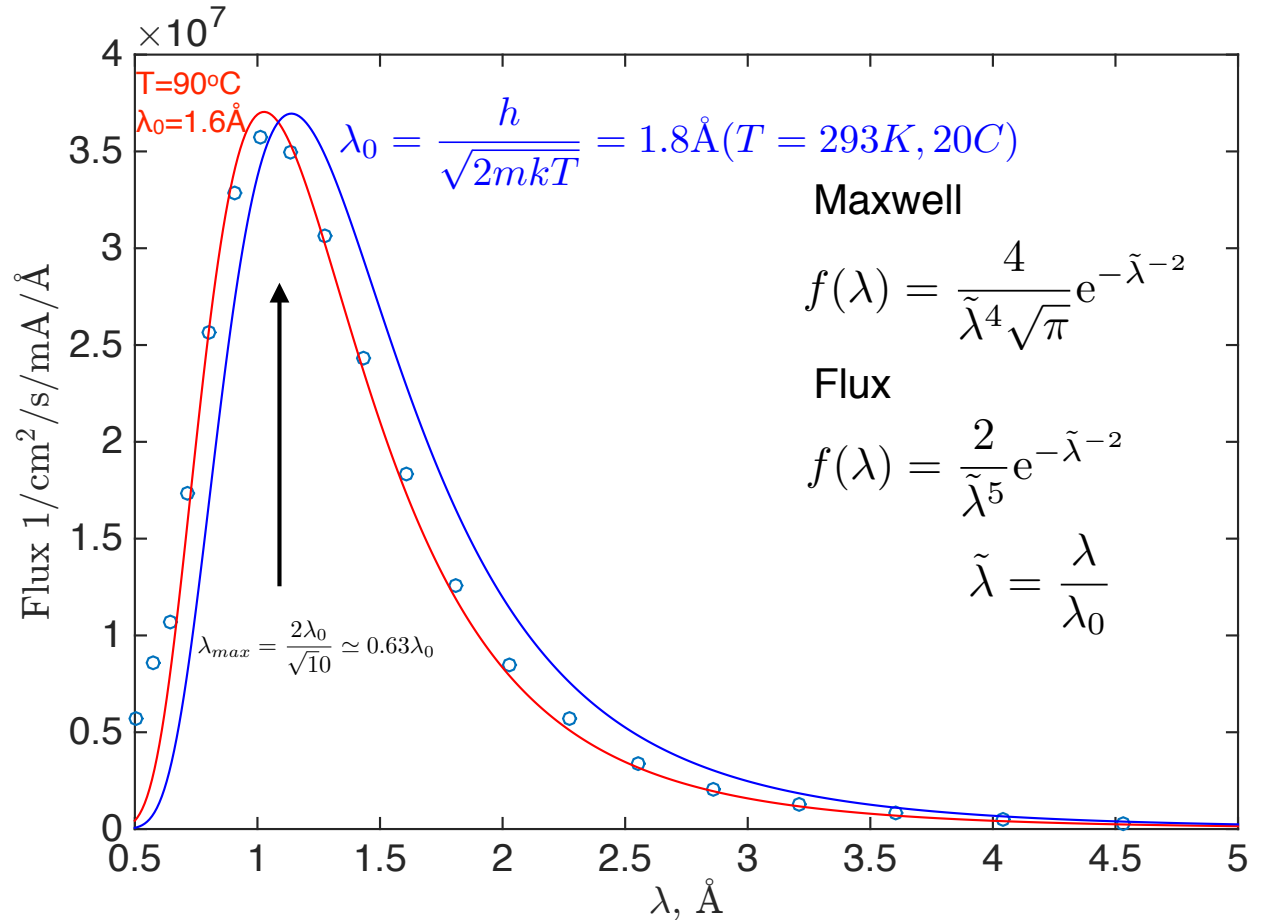
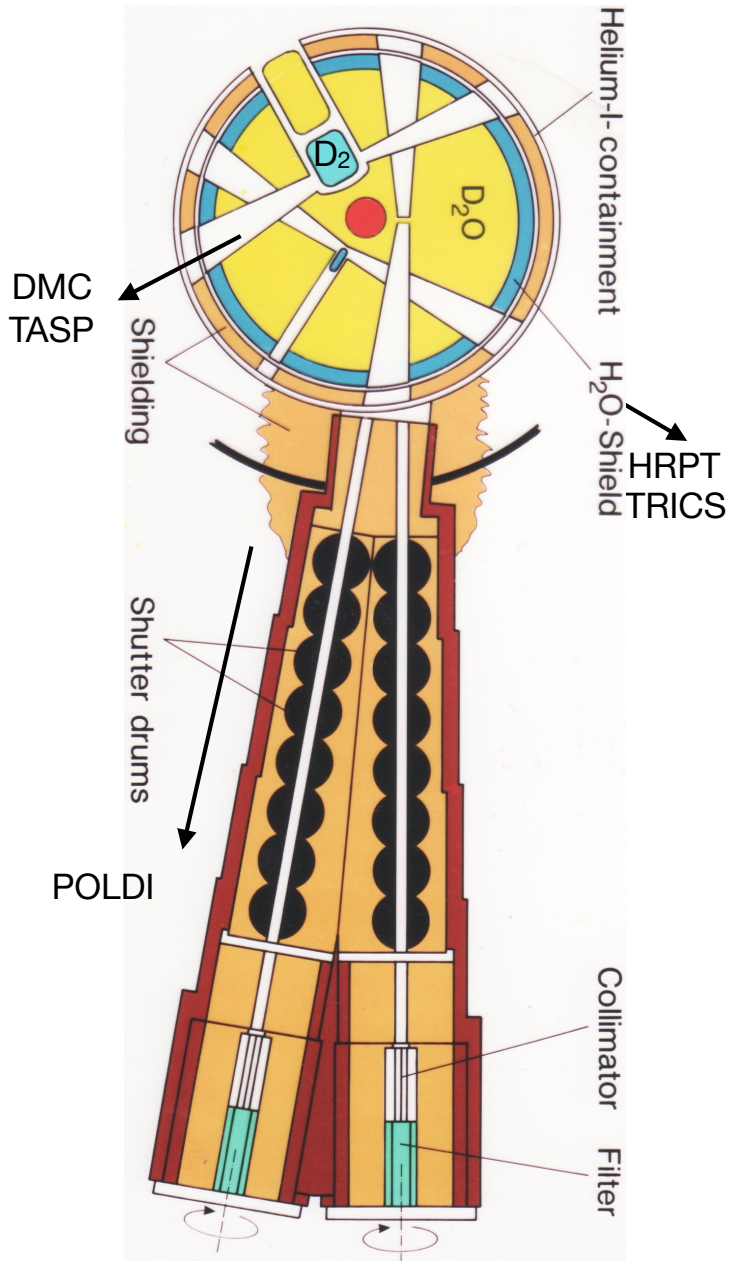
Calculated, from G.J. Russell, Spallation physics—an overview, Proceedings of ICANS-XI

SINQ hall



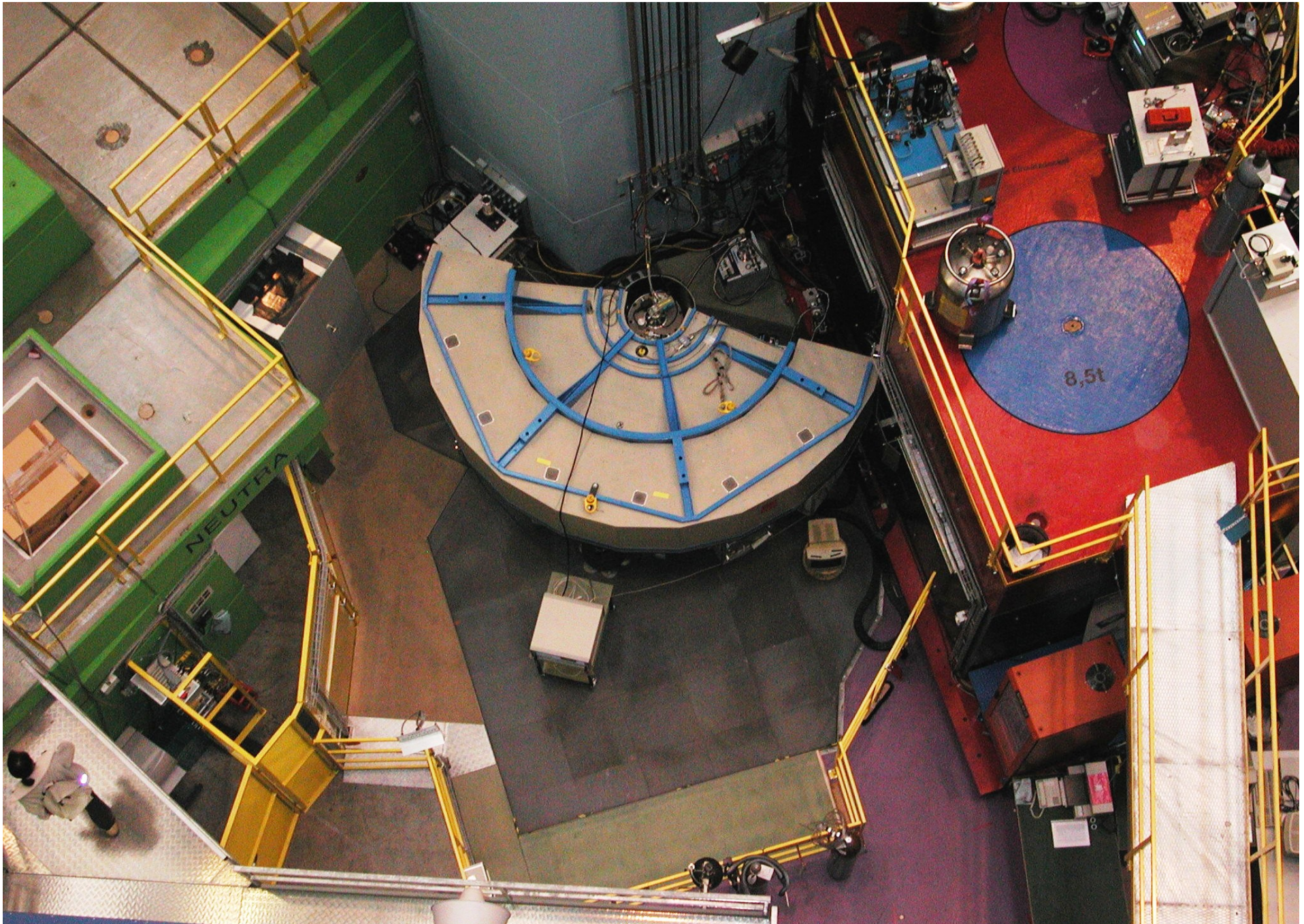
Elephant is:
Shielding of the direct
neutron beam also from
fast neutrons for
diffraction instruments

Neutron (thermal) flux from the D₂O moderator, Maxwellian at 90°C (HRPT, TRICS)



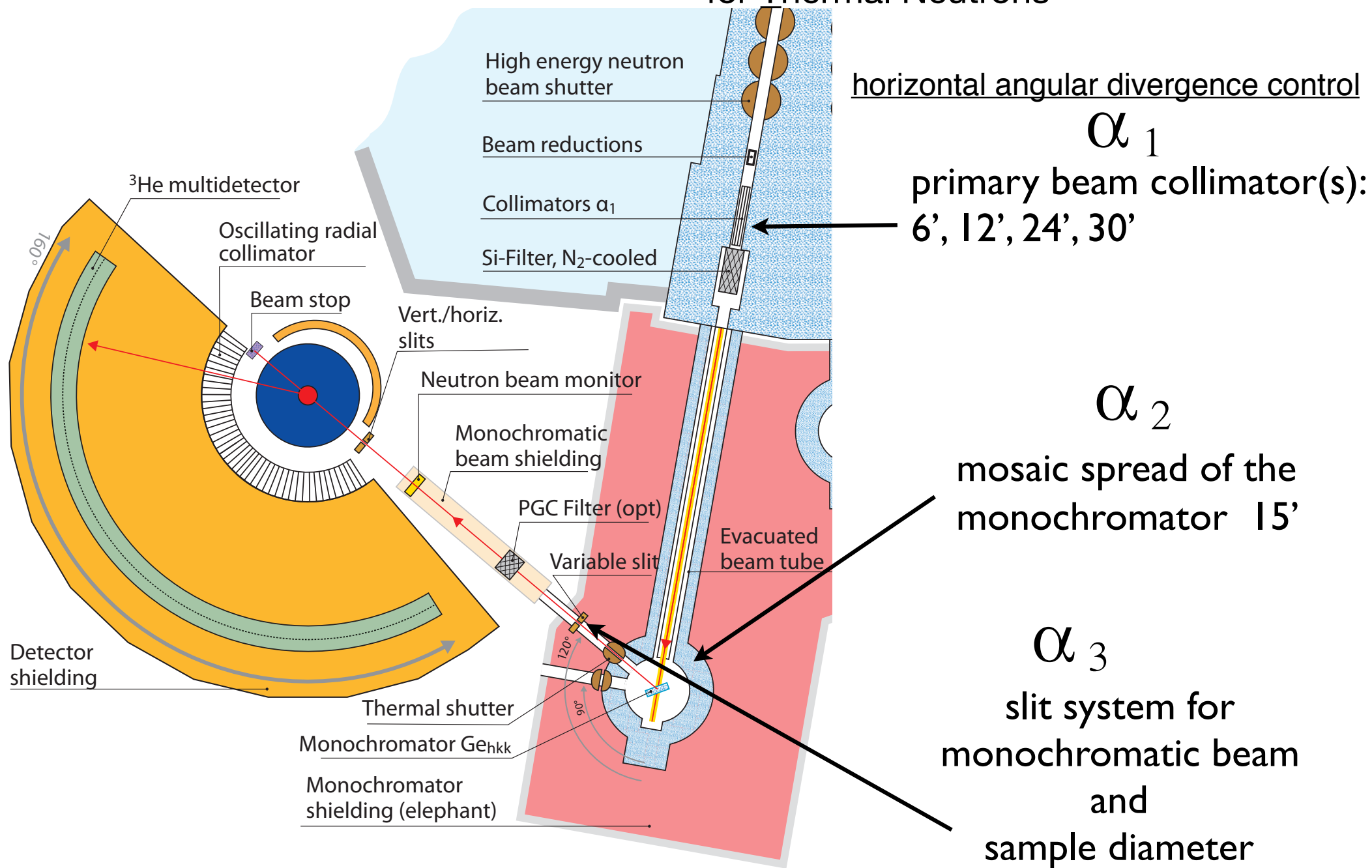
Total: $5 \cdot 10^7$ 1/cm²/s/mA
at SINQ current 2mA: 10^8

HRPT areal



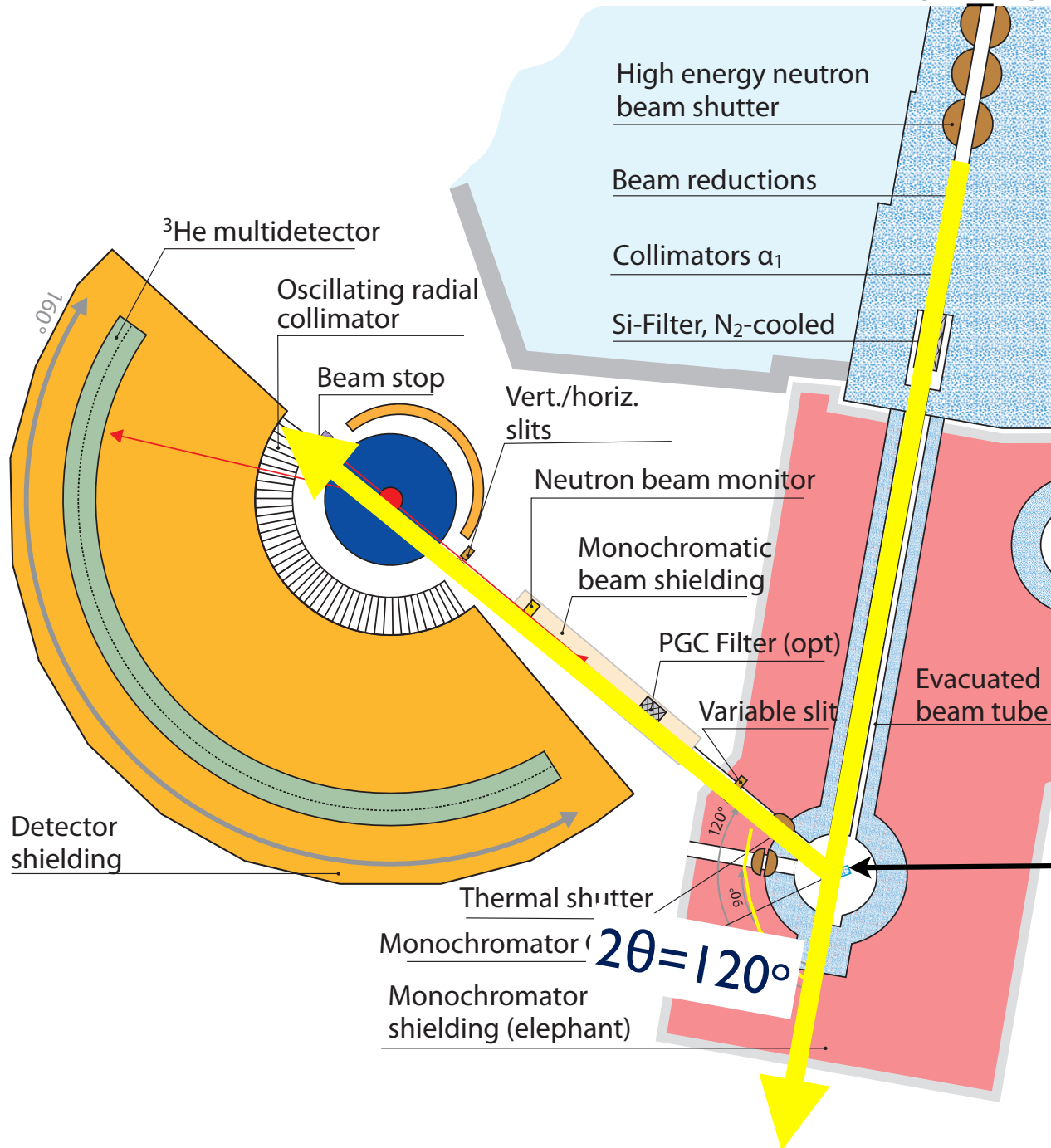
HRPT layout

High Resolution Powder Diffractometer for Thermal Neutrons



HRPT layout

High Resolution Powder Diffractometer for Thermal Neutrons

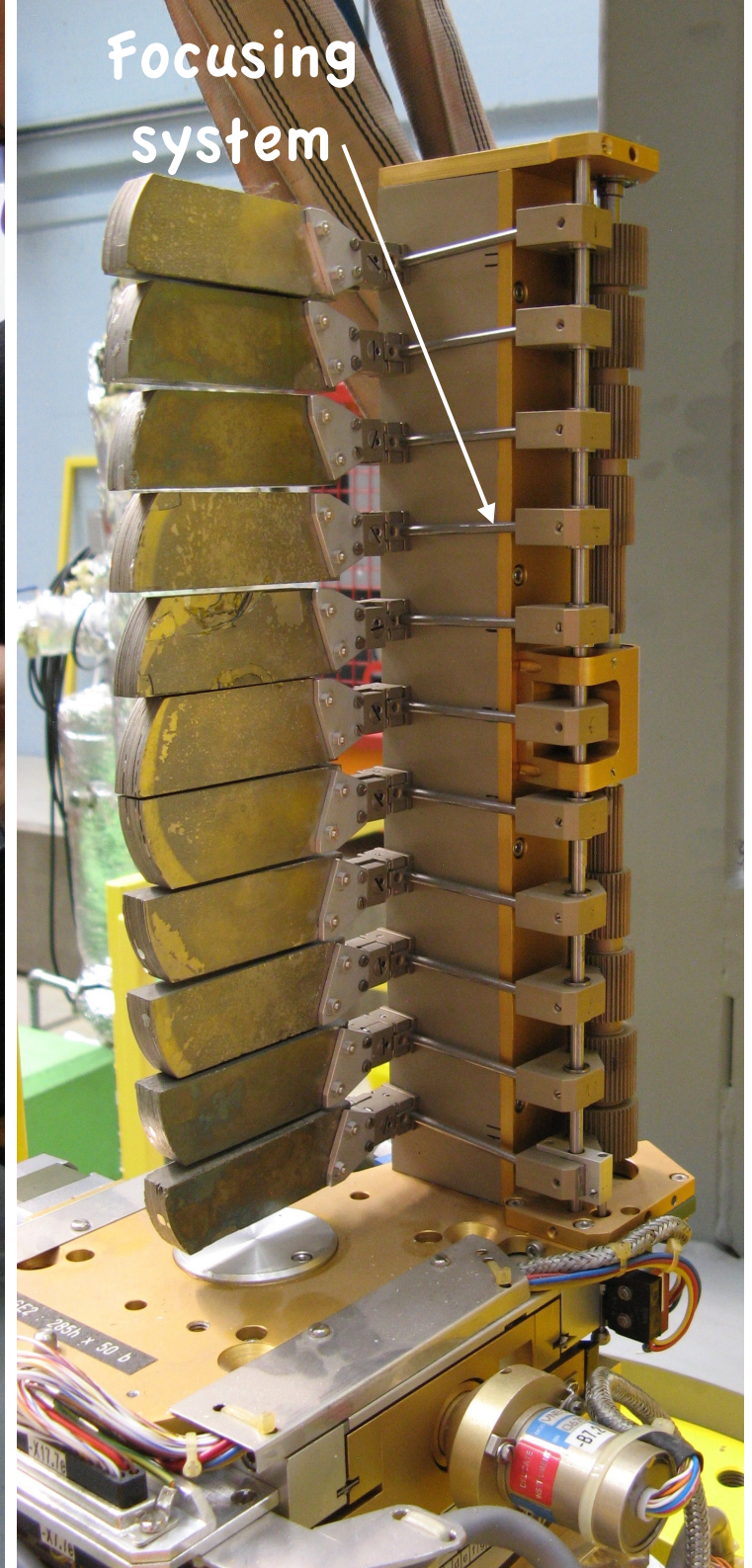


neutron monochromator
fixed 120 take off angle

$$\lambda = 2d \sin(\theta)$$

$$\lambda = 2d \sin(60^\circ)$$

Ge single crystal monochromator. 7 motors

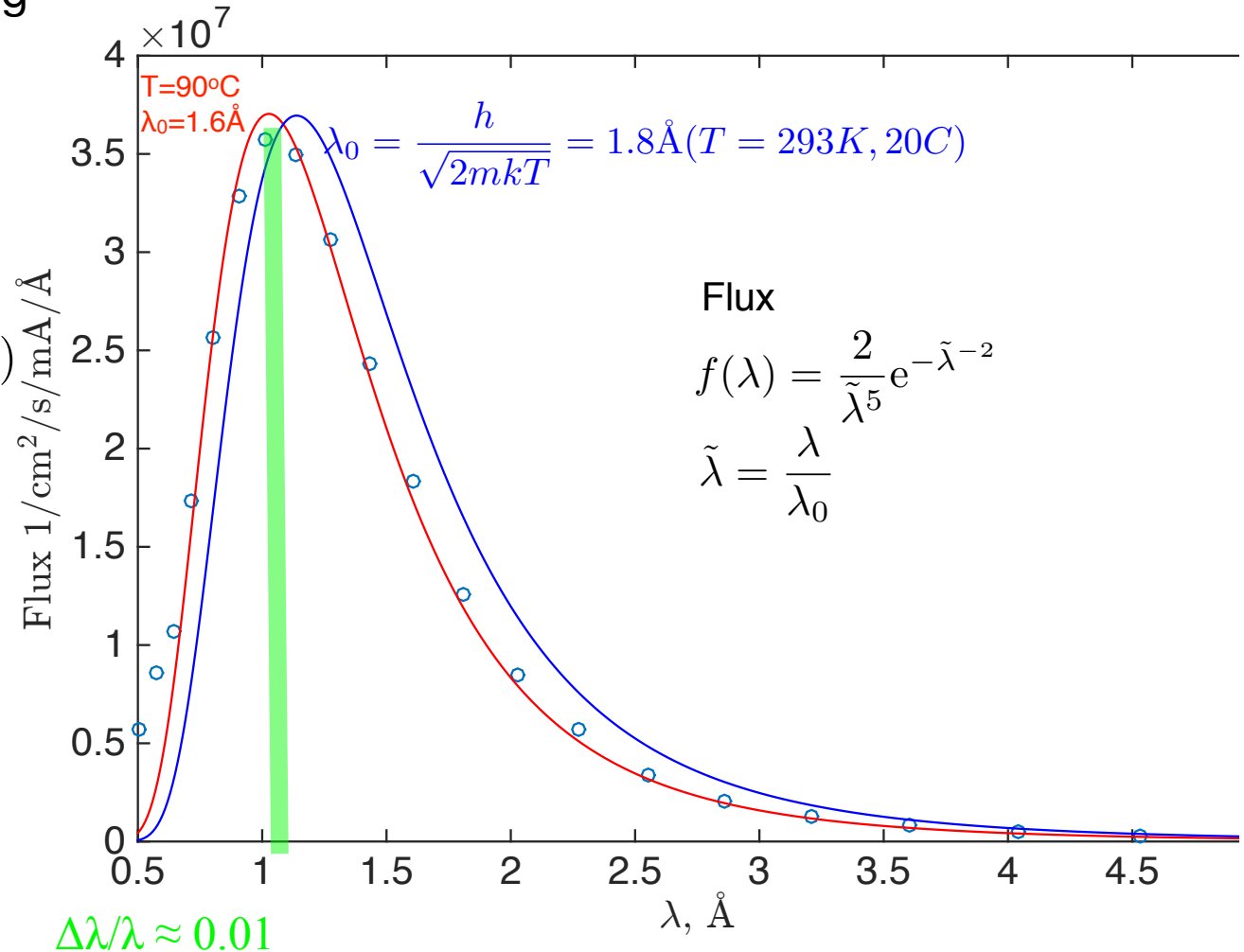


Monochromator cuts narrow wavelength range from the “white” flux. HRPT $\lambda=0.94 - 2.96 \text{ \AA}$

Intensity of Bragg scattering from big single crystal: Lorentz factor, extinction, geometry, ...

$$I \sim f(\lambda) \Delta\lambda C(\lambda, \theta) \sim f(\lambda) \lambda^{2.5} C'(\theta)$$

for fixed monochromator take-off 2θ for HRPT



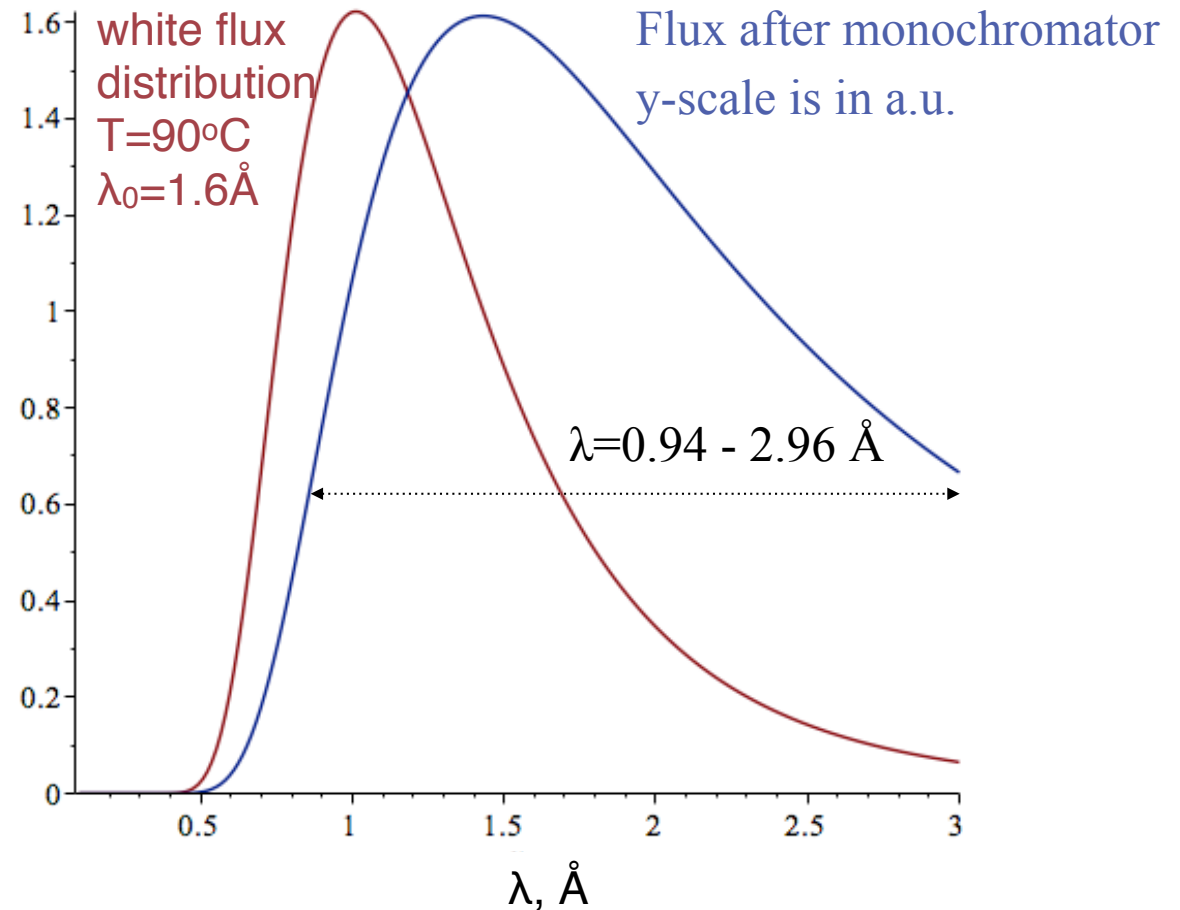
Total: $5 \cdot 10^7 \text{ 1/cm}^2/\text{s/mA}$
at SINQ current 2mA: 10^8

Monochromator cuts narrow wavelength range from the “white” flux. HRPT $\lambda=0.94 - 2.96 \text{ \AA}$

Intensity of Bragg scattering from big single crystal: Lorentz factor, extinction, geometry, ...

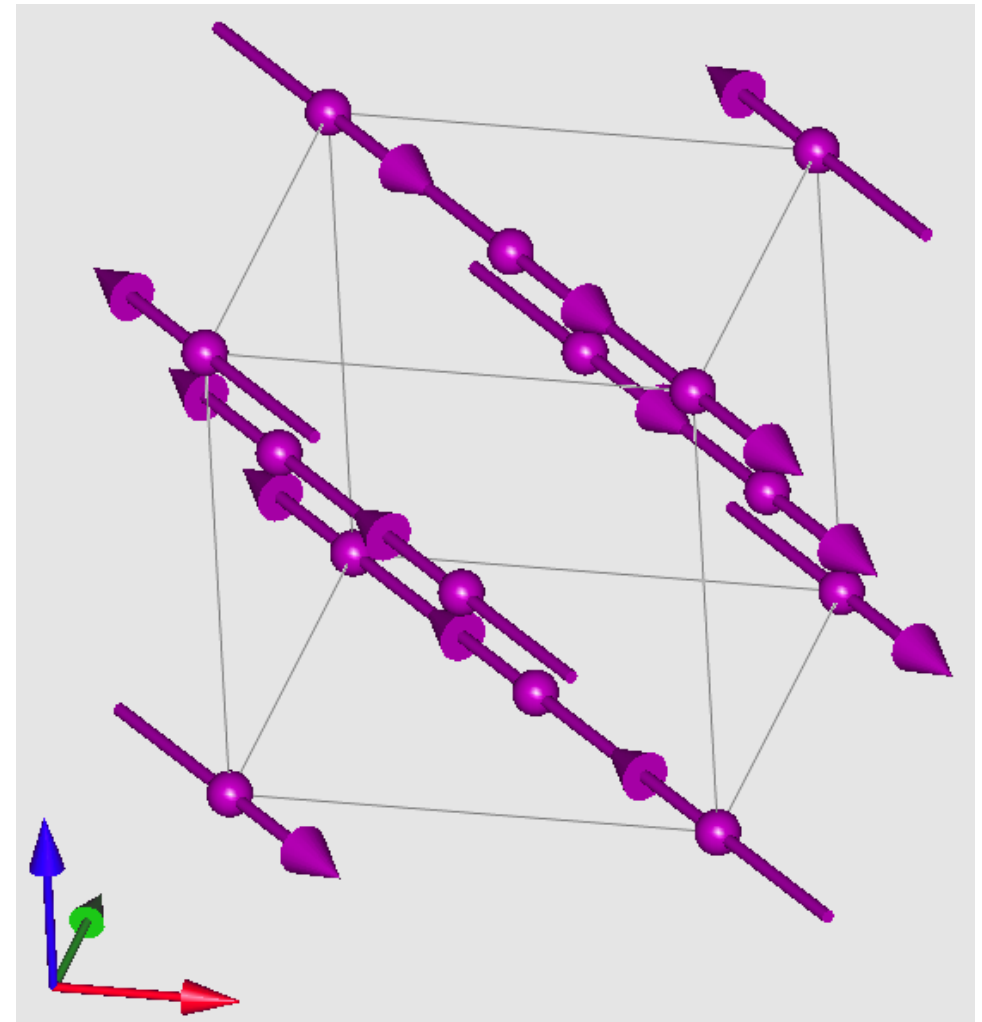
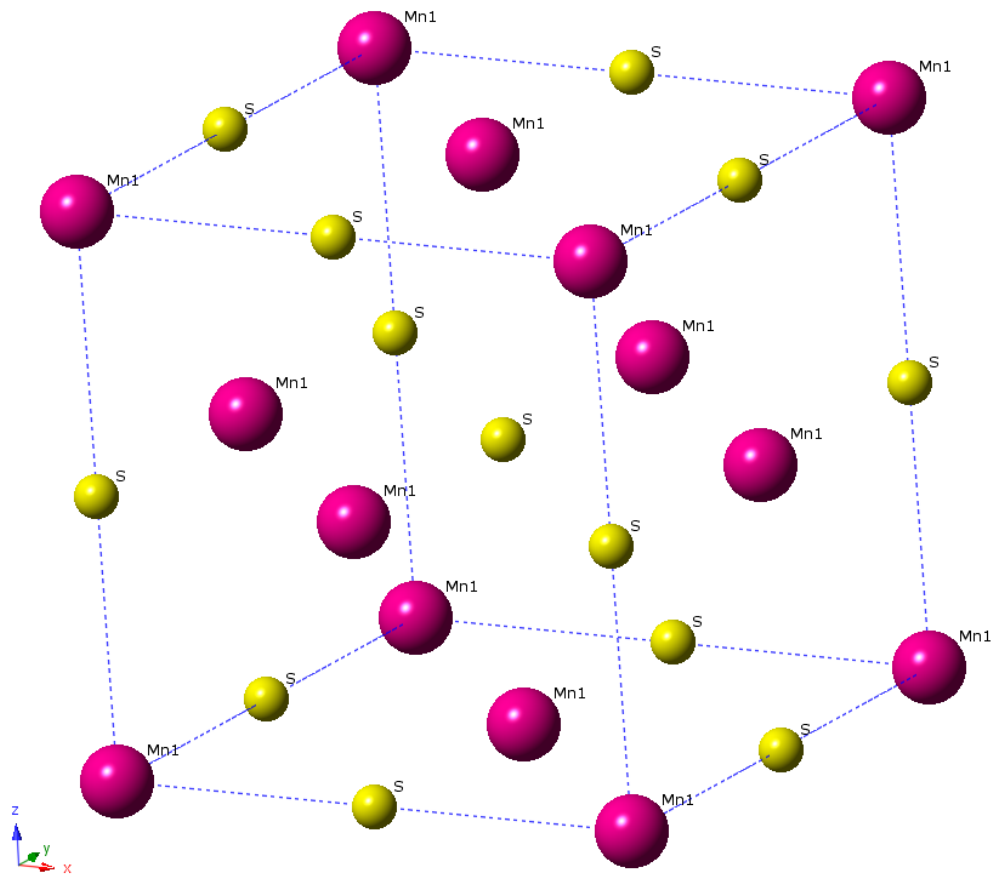
$$I \sim f(\lambda) \Delta\lambda C(\lambda, \theta) \sim f(\lambda) \lambda^{2.5} C'(\theta)$$

for fixed monochromator take-off 2θ for HRPT



Reminder on nuclear and magnetic neutron structure factors

Approximate crystal and magnetic structures of MnS below Néel temperature



neutron diffraction experiment ($\lambda = \text{const}$)

Nuclear structure factor

$$F(\mathbf{q}) = \sum_j b(\mathbf{r}_j) \exp(i\mathbf{q}\mathbf{r}_j)$$

Magnetic structure factor

$$\mathbf{F}(\mathbf{q}) \propto \sum_j \mathbf{S}_{0\perp j} \cdot \exp(i\mathbf{q}\mathbf{r}_j)$$

atom position and spin

Intensity in the detector

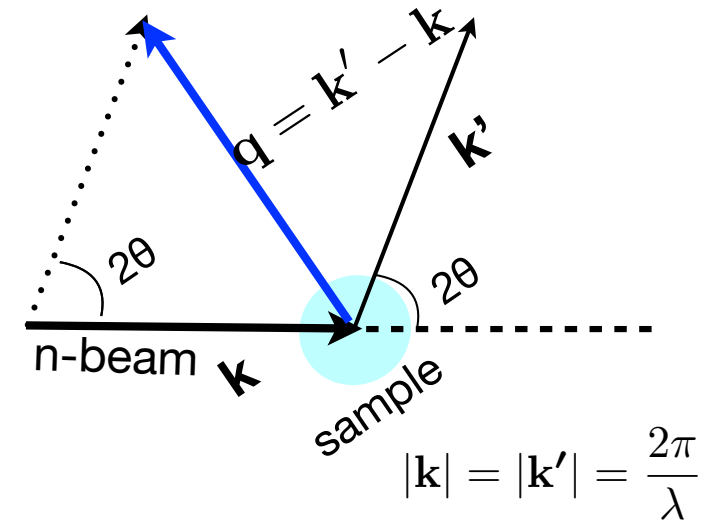
$$\frac{d\sigma}{d\Omega} \propto \mathbf{F}(\mathbf{q})\mathbf{F}^*(\mathbf{q}) \cdot \delta(\mathbf{H} - \mathbf{q})$$

momentum transfer or
scattering vector $\mathbf{q} = \mathbf{k}' - \mathbf{k}$

$$q = \frac{4\pi \sin(\theta)}{\lambda}, \theta = 0.. \leq 90\text{deg.}$$

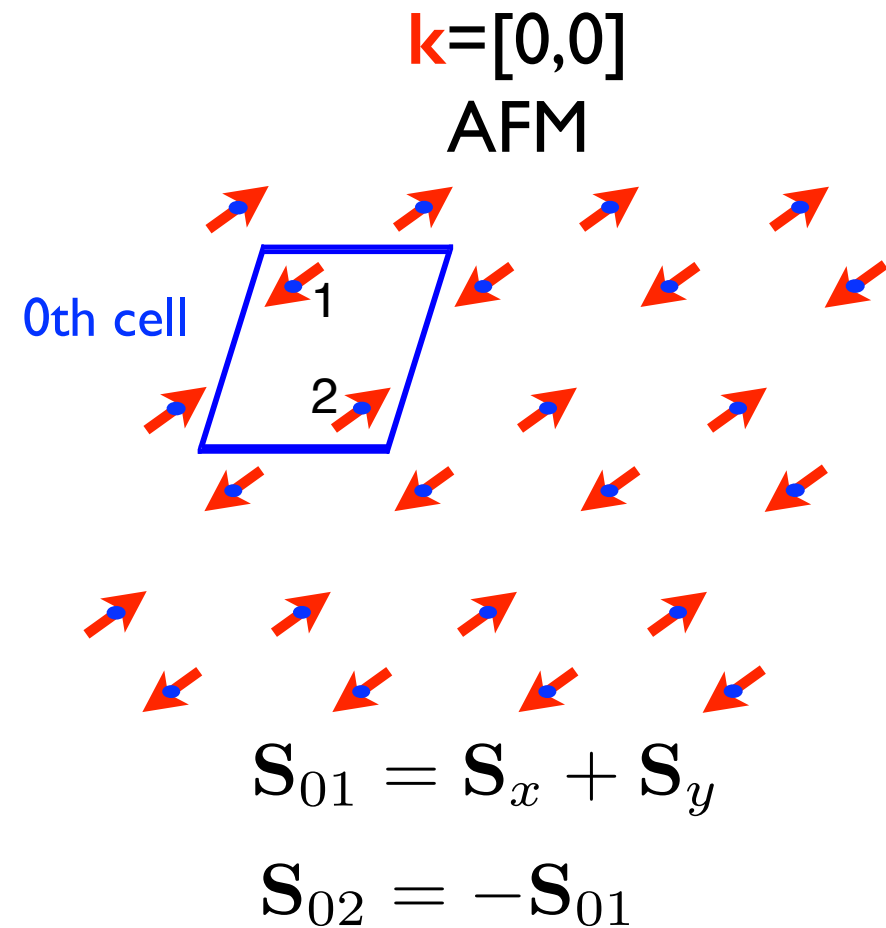
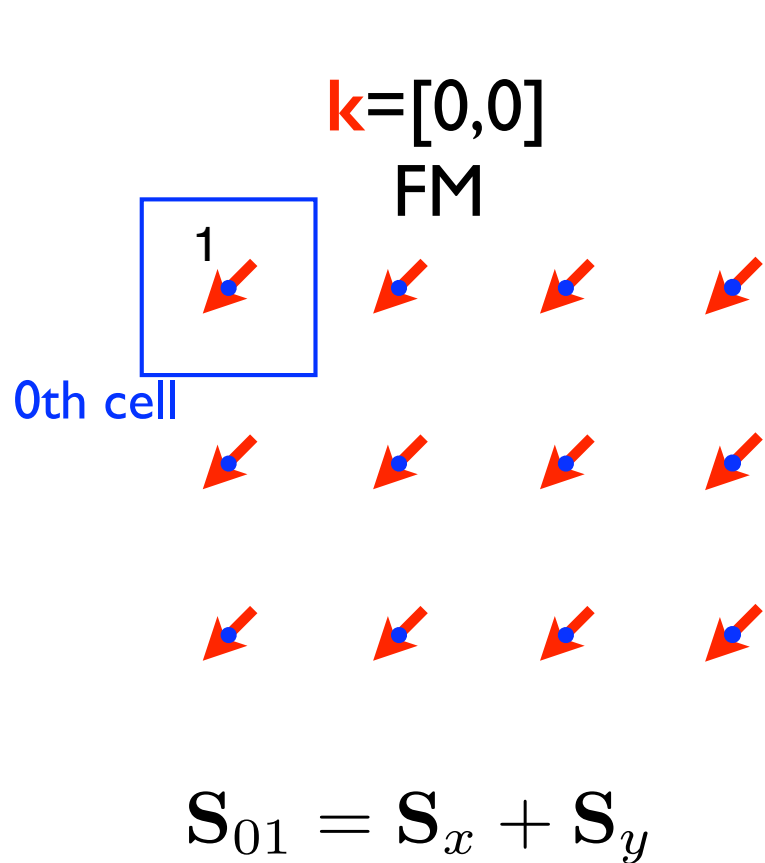
$$q_{\max} = \frac{4\pi}{\lambda} \quad \text{large } q_{\max} \rightarrow \text{good spatial resolution}$$

($d_{\min} = \frac{\lambda}{2}$)



Magnetic structure

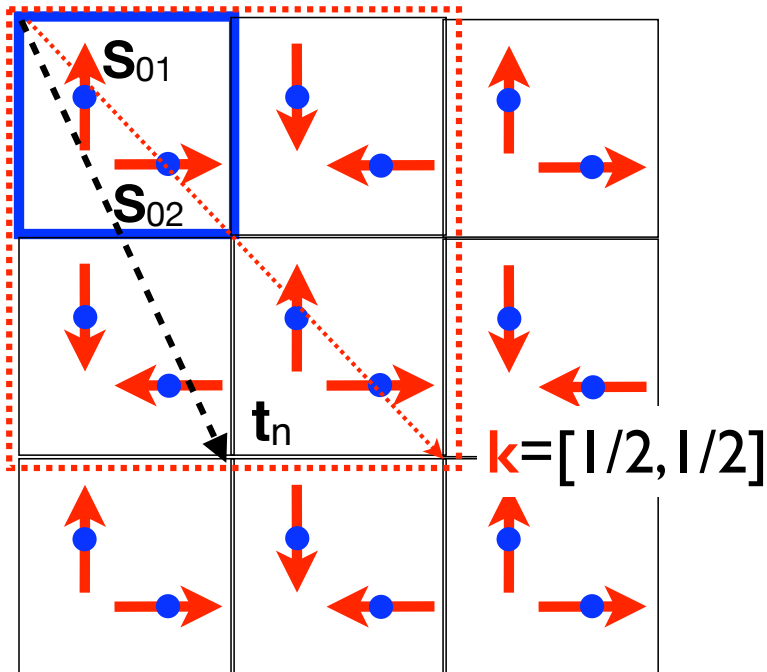
Examples



Examples of magnetic structures. Propagation vector formalism $\mathbf{k} \neq 0$.
 Magnetic mode $\mathbf{S}_0$ is specified in zeroth block of the cell == parent cell without centering translations

Magnetic moment
 below a phase transition

0th cell with many atoms in general



$$\mathbf{S}(\mathbf{t}_n) = \text{Re} \left(C \mathbf{S}_0 e^{2\pi i \mathbf{t}_n \mathbf{k}} \right) \sim \cos(2\pi \mathbf{t}_n \mathbf{k} + \varphi)$$

amplitude or
 mixing
 coefficients

magnetic mode

E.g., atom1 $\mathbf{S}_{01} = \mathbf{e}_y$
 atom2 $\mathbf{S}_{02} = \mathbf{e}_x$

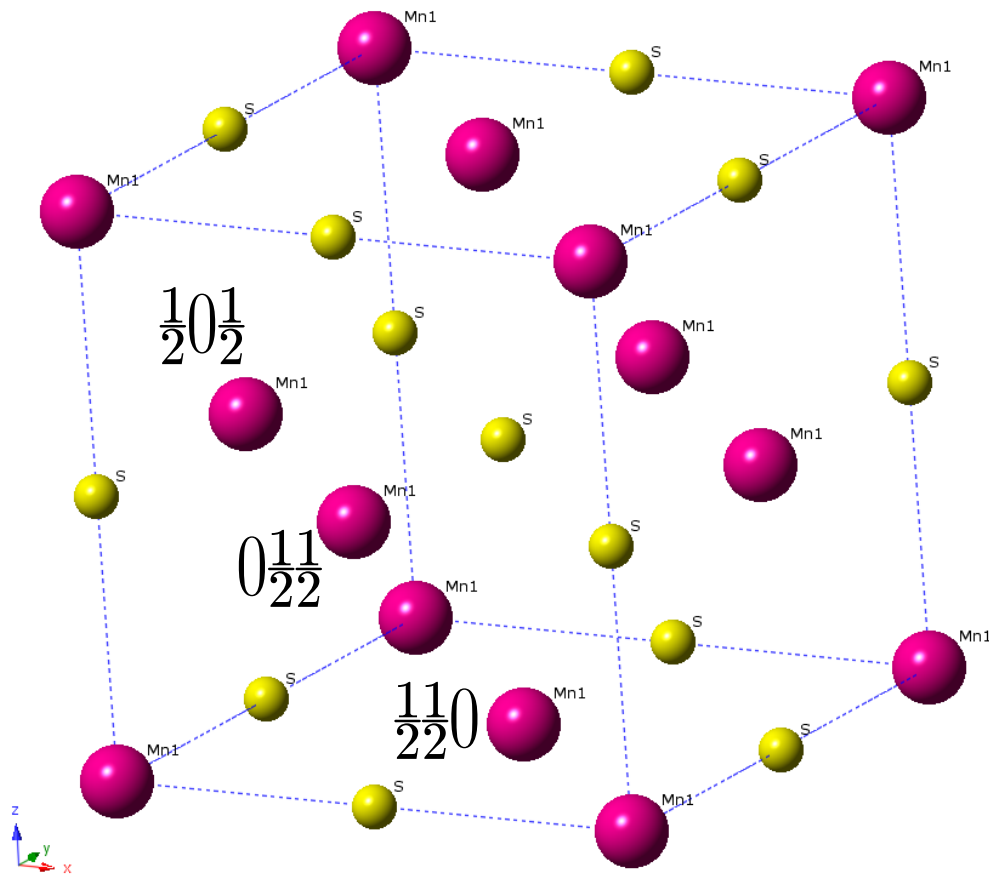
$$\mathbf{S}_1(\mathbf{t}_n) = C \mathbf{e}_y \cos(\pi(t_{nx} + t_{ny}))$$

$$\mathbf{S}_2(\mathbf{t}_n) = C \mathbf{e}_x \cos(\pi(t_{nx} + t_{ny}))$$

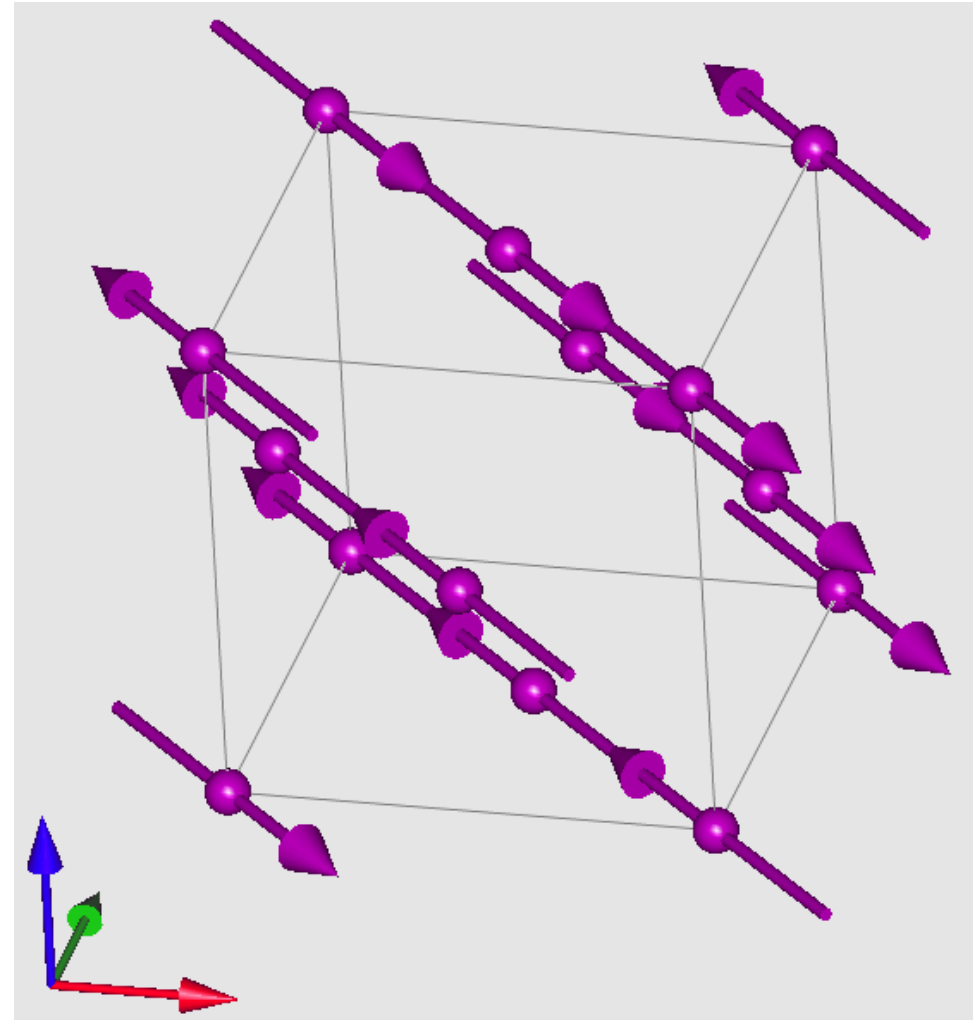
$$t_n = n \cdot 1 \text{ or } n \cdot \frac{1}{2}$$

Approximate crystal and magnetic structures of MnS below Néel temperature

single propagation vector
 $\mathbf{k} = [1/2, 1/2, 1/2]$



cubic, Fm-3m: Mn-atom in (000),
 three other Mn-atoms are generated
 by F-centering translations



Scattering from magnetic structure with propagation vector $\mathbf{k}$

In ND experiment we measure correlators of Fourier transform of magnetic lattice

$$\frac{d\sigma}{d\Omega} \propto (\mathbf{F}(\mathbf{q}) \cdot \mathbf{F}^*(\mathbf{q}) + i\mathbf{P} \cdot [\mathbf{F}(\mathbf{q}) \times \mathbf{F}^*(\mathbf{q})]) \cdot \delta(\mathbf{H} \pm \mathbf{k} - \mathbf{q})$$

structure factor

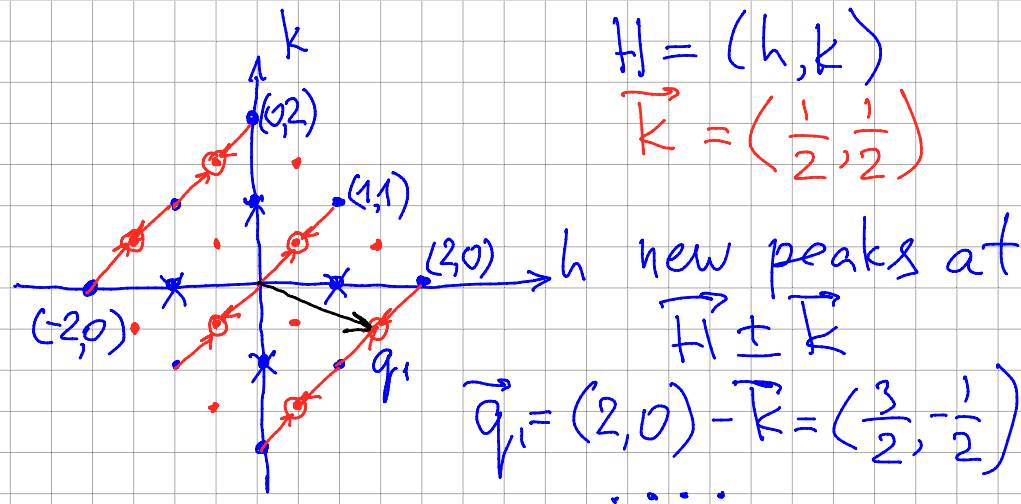
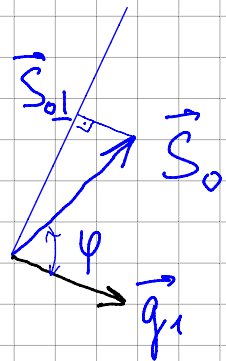
$$\mathbf{F}(\mathbf{q}) \propto \sum_j \mathbf{S}_{0\perp j} \cdot \exp(i\mathbf{q}\mathbf{r}_j)$$

polarized neutron
(chiral) term.

Bragg peak at
 $\mathbf{q} = \mathbf{H} \mp \mathbf{k}$

$$\vec{S}_\perp = \frac{\vec{q} \times \vec{S} \times \vec{q}}{|\vec{q}|^2}$$

$$|\vec{S}_\perp| = |\vec{S}_0| |\sin \varphi|$$



Practicum problems

MAGNETIC ORDER IN MnS

5. Practical course at SINQ

5.1 Manganese sulfide MnS

- rock salt crystal structure
- ionic crystal: Mn^{2+} , S^{2-}
- lattice constant $a = 5.199 \text{ \AA}$ at $T = 4.2 \text{ K}$
- space group $Fm\bar{3}m$
- electronic configuration of Mn^{2+} : $3d^5$
- Néel temperature $T_N = 161 \text{ K}$
- long-range antiferromagnetic order: antiferromagnetic stacking along (111) of ferromagnetic planes
- therefore doubling of the magnetic unit cell with respect to the crystallographic unit cell

Task 1: positions of nuclear Bragg peaks, indexing of the peaks

5.2 Neutron diffraction of MnS at room temperature

For all measurements at HRPT we will use $\lambda = 1.886 \text{ \AA}$

$$\lambda = 2d_{hkl} \sin \theta_{hkl}$$

Bragg law

(1)

λ : neutron wavelength, d_{hkl} : d-spacing of scattering plane hkl

θ_{hkl} : (half) scattering angle of reflection hkl in diffraction pattern

$$d_{hkl} = \frac{a}{\sqrt{h^2 + k^2 + l^2}}$$

(2)

a : cubic lattice constant, h, k, l indices of scattering plane

$$\vec{\tau}_{hkl} = \frac{2\pi}{a}(h, k, l)$$

$$\vec{\tau}_{hkl} \equiv \vec{H}$$

(3)

corresponding vector in reciprocal space: a node of reciprocal lattice

①

Tasks:

$$I(2\theta)$$

$$> T_N \approx 150 \text{ K}$$

- measure a diffraction pattern of MnS at $T = 300 \text{ K}$ in the paramagnetic state
- determine peak positions θ , d -spacings and indices (h, k, l) for all observed peaks

Coherent elastic cross section for nuclear neutron diffraction:

$$\frac{d\sigma}{d\Omega} \sim \sum_{\vec{\tau}_{hkl}} |F_{\vec{\tau}_{hkl}}|^2 \delta(\vec{Q} - \vec{\tau}_{hkl})$$

(4)

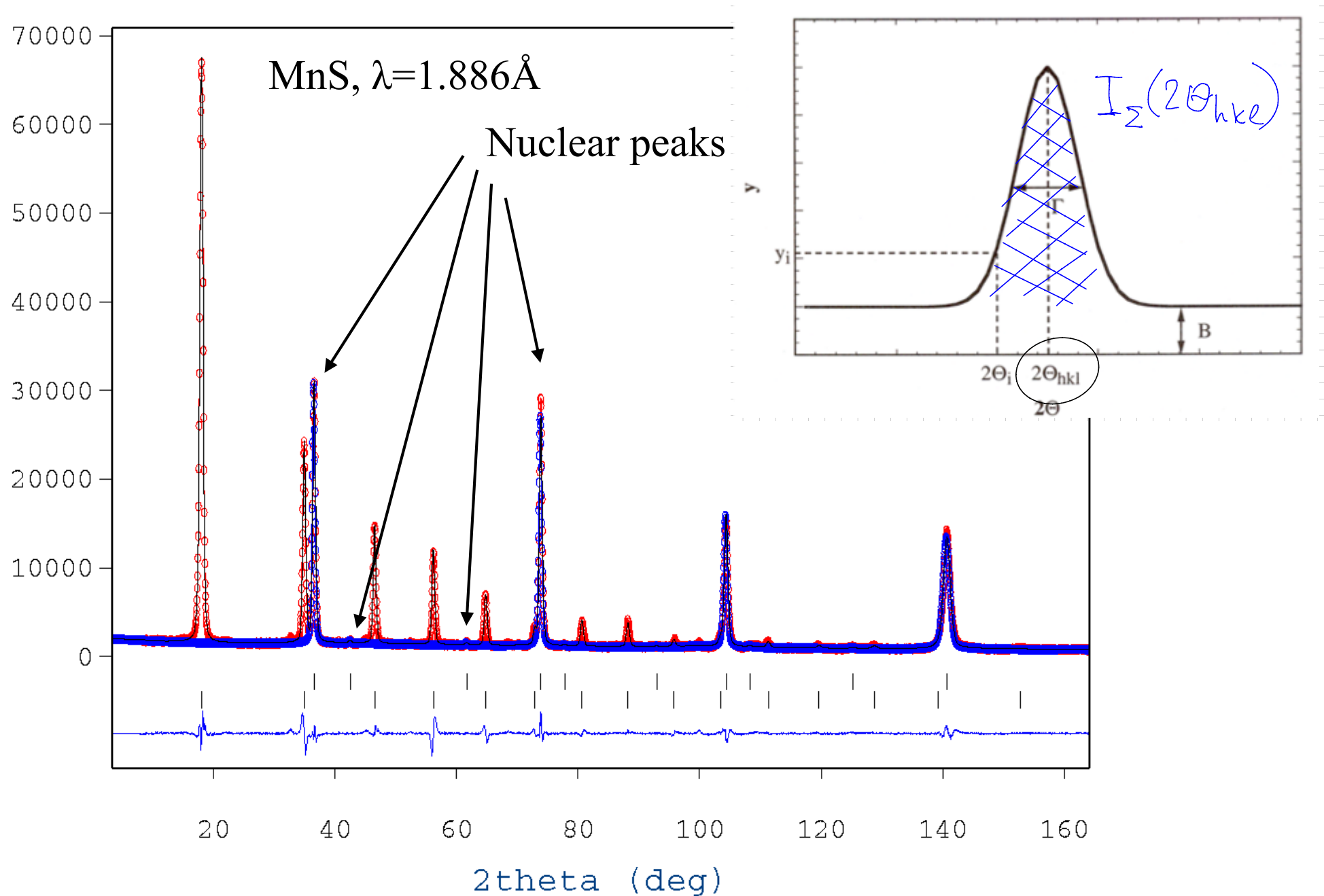
$$\vec{Q} = \vec{\tau}_{hkl}$$

$$|\vec{Q}| = \frac{4\pi \sin \theta}{\lambda}$$

$\vec{Q}$: scattering vector, $\vec{\tau}_{hkl}$: reciprocal lattice vector defining scattering planes

$\delta(\vec{Q} - \vec{\tau}_{hkl}) \rightarrow$ peak position given by crystal lattice (unit cell)

All the calculations/fits of experimental integrated intensities and peak positions will be done with 'fit' program under HRPT linux-computer



Task 2a: Calculation of structure factors and Bragg peak intensities and comparison with experiment

Coherent elastic cross section for nuclear neutron diffraction:

$$\frac{d\sigma}{d\Omega} \sim \sum_{\vec{\tau}_{hkl}} |F_{\vec{\tau}_{hkl}}|^2 \delta(\vec{Q} - \vec{\tau}_{hkl}) \quad (4) \quad \vec{Q} = \vec{\tau}_{hkl} \quad |\vec{Q}| = \frac{4\pi \sin \theta}{\lambda}$$

$\vec{Q}$: scattering vector, $\vec{\tau}_{hkl}$: reciprocal lattice vector defining scattering planes

$\delta(\vec{Q} - \vec{\tau}_{hkl}) \rightarrow$ **peak position given by crystal lattice (unit cell)**

$F_{\vec{\tau}_{hkl}}$: structure factor of unit cell

$$F_{\vec{\tau}_{hkl}} = \sum_{\vec{d}_i} \left(b_{\vec{d}_i} e^{i\vec{\tau}_{hkl} \cdot \vec{d}_i} \cdot e^{-B_i (Q/4\pi)^2} \right), \text{ we assume } B=0, \text{ but actually } B \approx 0.8 \text{ \AA}^2 \text{ at RT} \quad (5)$$

$\vec{d}_i$: atomic coordinate of i-th atom in real space, sum runs over all atoms in unit cell

$b_{\vec{d}_i}$: scattering length of atom at position $\vec{d}_i$

Intensity $\sim |F_{\vec{\tau}_{hkl}}|^2 \rightarrow$ **peak intensity is mainly given by arrangement of atoms in unit cell**

The sum runs over all atoms in unit cell.

For MnS:

$$b_{Mn} = -3.73 \text{ fm}, \quad b_S = 2.85 \text{ fm}$$

femto or fermi
 $1 \text{ fm} = 10^{-13} \text{ cm} = 10^{-15} \text{ m}$

$\vec{d}$ -vectors:

Mn:	$\vec{d}_1 = a(0, 0, 0)$	$\vec{d}_2 = a(0, \frac{1}{2}, \frac{1}{2})$
	$\vec{d}_3 = a(\frac{1}{2}, 0, \frac{1}{2})$	$\vec{d}_4 = a(\frac{1}{2}, \frac{1}{2}, 0)$
S:	$\vec{d}_5 = a(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$	$\vec{d}_6 = a(\frac{1}{2}, 0, 0)$
	$\vec{d}_7 = a(0, \frac{1}{2}, 0)$	$\vec{d}_8 = a(0, 0, \frac{1}{2})$

F-centering. (8)

Task 2b: Calculation of structure factors and Bragg peak intensities and comparison with experiment

$I_{\Sigma}(2\theta)$ is measured in the experiment $|\vec{Q}| = \frac{4\pi \sin \theta}{\lambda}$

For cylindrical geometry of the powder sample container the integrated intensity of the scattered neutrons of the Bragg peak at $|\vec{Q}|$ is given by

$$I(Q) = C \cdot A(\theta) \cdot L(\theta) \cdot \frac{d\sigma}{d\Omega} = C \cdot A(\theta) \cdot L(\theta) \cdot |F(Q)|^2 \cdot \text{mult}$$

assume: $A(\theta) \approx 1$ (6)

C: scale factor, $A(\theta)$: absorption factor, $L(\theta)$: Lorentz factor, *mult*: multiplicity

$$L(\theta) = \frac{1}{\sin \theta \sin 2\theta}$$

(7)

The Lorentz factor $L(\theta)$ is a geometrical correction depending on the scattering geometry.

2

Tasks:

- calculate $|F_{\vec{\tau}_{hkl}}|^2$ for $(h, k, l) = (1, 1, 1)$, calculate the multiplicity
- calculate $|F_{\vec{\tau}_{hkl}}|^2$ for $(h, k, l) = (2, 0, 0)$, calculate the multiplicity
- compare ratio of $I(\vec{\tau}_{111}) / I(\vec{\tau}_{200})$ with the measured intensity ratio

$\rightarrow (111), (-111), \dots$

Task 3: Indexing of the magnetic Bragg peaks. Calculation of magnetic structure factors and determination of the value and direction of the Mn-spins.

5.3 Neutron diffraction of MnS in the magnetically ordered state

3

Tasks:

- measure a neutron diffraction pattern of MnS at $T = 50$ K in the magnetically ordered state or lower T
- compare this data with the paramagnetic pattern at room temperature
- index the magnetic peaks, i.e. find (h, k, l) for ^{first 3} each magnetic peak
- based on the indices, what is the magnetic unit cell compared to the crystallographic one

- using eq. (6) for (111) -peak and eq. (10) for $(\frac{1}{2}\frac{1}{2}\frac{1}{2})$ magnetic peak determine $\vec{\mu}$.

Consider $\vec{\mu} \parallel (111)$ in calculations

- Which direction of $\vec{\mu}$ can be excluded?

Task 3: Indexing of the magnetic Bragg peaks. Calculation of magnetic structure factors and determination of the value and direction of the Mn-spins.

Coherent elastic cross section for

antiferromagnetic order

$$\frac{d\sigma}{d\Omega} \sim \sum_{\vec{r}_{M,hkl}} |F_{M,hkl}|^2 \delta(\vec{Q} - \vec{r}_{M,hkl})$$

$F_{M,hkl}$: antiferromagnetic structure factor

Formula (10) is actually formula (6) with $|F_{M\perp}|$ instead of $|F|^2$ with the same scale factor C

The intensity of the magnetic Bragg peak at $|\vec{Q}_M|$ is

$$I(Q_M) = C \cdot A(\theta) \cdot L(\theta) \cdot |F_{M\perp}|^2 \cdot \text{mult}$$

where

$$\vec{F}_{M\perp} = \frac{1}{2} r_0 \sum_{j=1}^4 e^{i\vec{Q}_M \cdot \vec{r}_j} \vec{\mu}_{j\perp} \quad (11)$$

$$\text{and } \vec{\mu}_{\perp} = \left(\vec{\mu} - \frac{\vec{Q}_M (\vec{\mu} \cdot \vec{Q}_M)}{Q_M^2} \right) \quad (12)$$

where $\vec{\mu}$ is the magnetic moment in units μ_B and $r_0 = -0.54 \cdot 10^{-12}$ cm, $\vec{Q}_M \equiv \vec{r}_{M,hkl}$

The sum runs over Mn-atoms.

$\vec{\mu}_1 \uparrow \downarrow \vec{\mu}_{2,3,4}$ Can you explain why?